### A planar-algebraic universe

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The University of Queensland

#### Aalto University Mathematical Physics Seminar

Some suggestions that spacetime is not a continuum:

- Aspects of the 'It from Qubit' program
- Successes of loop quantum gravity and causal sets
- Computational universe hypothesis

#### Question

Can we discretise our theories such that they are recovered under a continuum limit?

### Outline



- 2 Discrete conformal nets
- Semicontinuous models
- Integrable operators

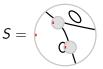


# Planar algebras

### Planar algebras

#### **Planar tangles**



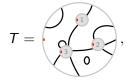


Graded vector space

 $(P_n)_{n\in\mathbb{N}_0}$ 

Multilinear maps

 $\mathsf{P}_T \colon P_2 \times P_4 \times P_6 \to P_8$ 



$$\mathsf{P}_{T}(v_{1}, v_{2}, v_{3}) = \bigcirc_{v_{3}}^{(v_{1})} 0 \in P_{8}$$

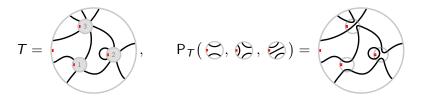
### Example: Temperley-Lieb planar algebra

Vector spaces

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Planar tangle action

 $P_T: P_4 \times P_6 \times P_6 \rightarrow P_8$ 



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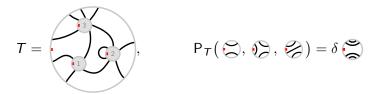
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Planar tangle action

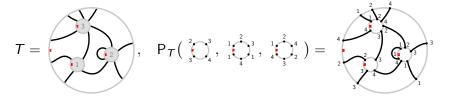
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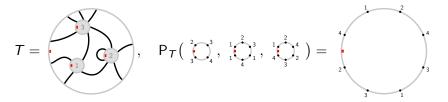


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**Planar tangle action** 

$$T = \bigvee_{1}^{13} (1 + 2)^{13} ($$

### Annular tangles

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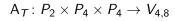


Graded annular vector space

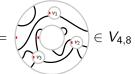
$$(V_{m,n})_{m,n\in\mathbb{N}_0}$$



#### Multilinear maps



$$\mathsf{A}_{T}(v_1,v_2,v_3) =$$





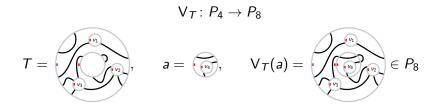
Xavier Poncini

### Action of annular vectors

#### Annular vectors



#### Linear maps

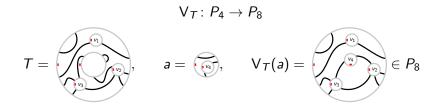


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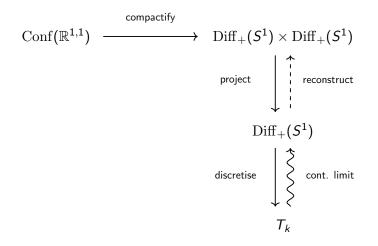
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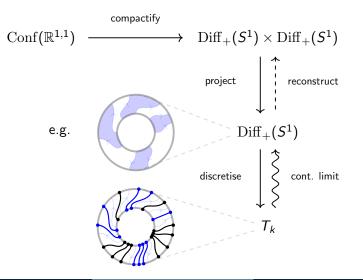
#### Linear maps



# 'Conformal' groups



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A conformal net (Haag-Kastler '64) consists of:

- i) a Hilbert space  ${\cal H}$
- ii) a  $\mathit{C}^*\text{-algebra}\ \mathcal{A}(\mathit{I})$  on  $\mathcal{H}$  for each open interval  $\mathit{I}\subset S^1$
- iii) a continuous unitary representation U of  $\mathrm{Diff}_+(S^1)$  on  $\mathcal H$

Subject to:

Isotony:  $\mathcal{A}(I) \subseteq \mathcal{A}(J)$  if  $I \subseteq J$ Locality:  $[\mathcal{A}(I), \mathcal{A}(J)] = 0$  if  $I \cap J = \emptyset$ Covariance:  $U(\alpha)\mathcal{A}(I)U(\alpha)^* = \mathcal{A}(\alpha(I))$   $\alpha \in \text{Diff}_+(S^1)$ 

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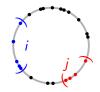


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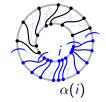


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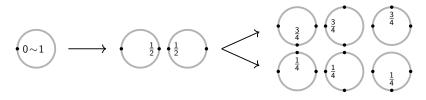
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 $\mathcal D$  is the set of k-adic subdivisions of  $S^1$  for  $k\in\mathbb N_{\ge2}$ 



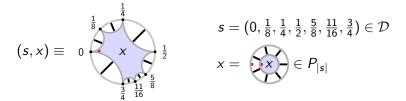
where  $S^1$  is divided into intervals of the form  $[\frac{m}{k^n}, \frac{m+1}{k^n}]$  for  $m, n \in \mathbb{N}_0$ .

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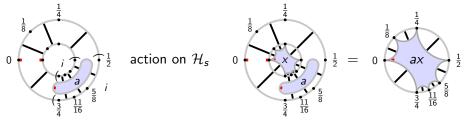
vectors in  $\mathcal{H}_s$  are of the form:



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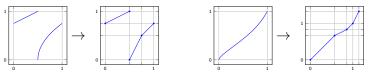
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### D is Thompson's group $T_k$



### Theorem (Zhuang '07)

For each  $f \in \text{Diff}_+(S^1)$  there exists a  $g \in T_k$  approx. f to arb. precision.

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Unitary representation induced by a map

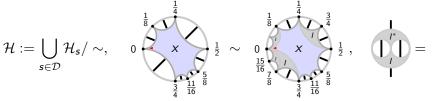
$$\Phi: T_k \to (V_{m,n})_{m,n \in \mathbb{N}_0}$$

Continuity: work in progress!

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The full Hilbert space  $\mathcal{H}$  is defined



where  $R \in P_{k+1}$ . The algebras  $\mathcal{A}(i)$  are defined similarly.

# Continuous representations

Definition

r

A representation  $\pi$  is *continuous* if each sequence  $(f_n)_{n \in \mathbb{N}} \subset T_k$  satisfies

$$\lim_{n\to\infty} \|f_n - \mathrm{id}\| = 0, \qquad \lim_{n\to\infty} \langle x, \pi(f_n)(y) \rangle = \langle x, y \rangle, \qquad \forall \, x, y \in \mathcal{H}.$$

Denote by  $\operatorname{Rot}_k$  the rotation subgroup of  $T_k$ , generated by:

$$\varrho_s: S^1 \to S^1, \qquad \qquad x \mapsto x + s \mod 1,$$

where s is a k-adic rational. Matrix elements can be expressed as:

$$\langle x, y \rangle = \underbrace{\bigvee_{x}^{y^*}}_{x}, \qquad \langle x, U_l(\varrho_{\frac{1}{k^r}})(y) \rangle = \underbrace{\bigvee_{x}^{y^*}}_{x}$$

where  $x, y \in \mathcal{H}$ .

Xavier Poncini

# Continuity conditions

For each  $x, y \in \mathcal{H}$ , there exists a sufficiently large  $r \in \mathbb{N}$ , such that

$$\langle x, U_{l}(\varrho_{\frac{1}{k^{r}}})(y) \rangle = \bigvee_{x}^{y^{*}} = \bigvee_{x}^{r} \bigvee_{$$

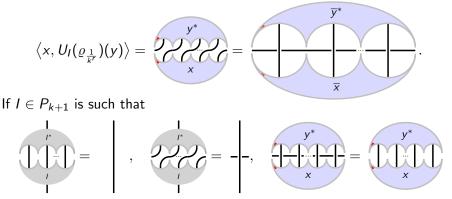
for all  $x, y \in \mathcal{H}$ . Then

$$\lim_{r\to\infty} \langle x, U_l(\varrho_{\frac{1}{k^r}})(y) \rangle = \langle x, y \rangle,$$

and with other arguments  $U_I : \operatorname{Rot}_k \to \operatorname{U}(\mathcal{H})$  is a continuous unitary rep.

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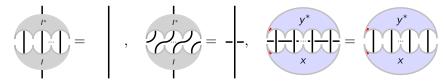
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#### If $I \in P_{k+1}$ is such that



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and with other arguments  $U_I : \operatorname{Rot}_k \to \operatorname{U}(\mathcal{H})$  is a continuous unitary rep.

Xavier Poncini

#### Brauer algebra solution

The Brauer planar algebra  $(P_n)_{n \in 2\mathbb{N}_0}$  is generated by the action of planar tangles on the space  $P_4 = \operatorname{span}(\{ \bigoplus , \bigoplus , \bigoplus \})$ , subject to:

$$\bigotimes = \bigotimes \qquad \bigotimes = \delta \qquad \bigotimes = \delta \qquad (\bigotimes = \delta)$$
  
Specialising  $k = 5$  and  $\delta = 1$  we have the solution  $\bigvee = \bigvee$ 

This solution can be generalised to k = 2n + 5 for all  $n \in \mathbb{N}_0$ 

$$I = \bigvee_{n} P_{2n}.$$

#### Theorem

For  $I \in P_{2n+6}$  above  $U_I$  is a continuous unitary representation of  $\operatorname{Rot}_{2n+5}$ .

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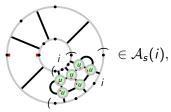
# Integrable operators

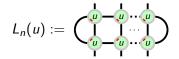
### Spin chains on spacetime

*R*-operators:  $u \in \mathbb{C}$ 

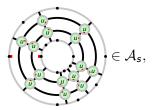
$$R(u) = \mathbf{y} = \sum_{a \in B_A} r_a(u) a$$

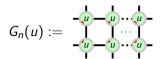
Local transfer operators





**Global transfer operators** 



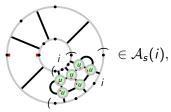


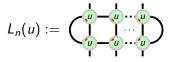
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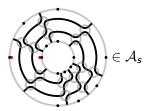
$$R(u) = \sum_{a \in B_a} r_a(u) a$$

Local transfer operators





Global transfer operators: e.g spacetime translations are generated by



### Integrability

A transfer operator  $T(u) \in A_s(i)$  is **integrable** if

$$[T(u), T(v)] = 0, \qquad \forall u, v \in \Omega \subseteq \mathbb{C}.$$

Expanding T(u) in a basis of scalar functions

$$\mathcal{T}(u) = \sum_{i=0}^{\infty} u^i Q_i$$
, integrability implies  $[Q_i, Q_j] = 0$ ,  $\forall i, j \in \mathbb{N}_0$ ,

where  $H \equiv Q_1$  is the hamiltonian. T(u) is **polynomially integrable** if

 $\exists \; b \in \mathcal{A}_{s}(i) \; \; ext{such that} \; \mathcal{T}(u) \in \mathbb{C}(u)[b]$ 

#### Theorem (XP, Rasmussen '22)

If T(u) is integrable and diagonalisable it is polynomially integrable.

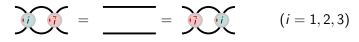
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## Yang-Baxter integrability

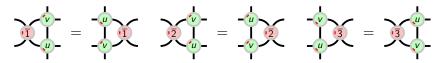
A model is **Yang-Baxter integrable** if the *R*-operators satisfy *local* relations that imply [T(u), T(v)] = 0.

For  $T(u) \in \{L_n(u), G_n(u)\}$  a set of sufficient conditions is given by:

Inversion identities



• Yang-Baxter equations



• Boundary Yang-Baxter equations



## Polynomial integrability

There is an adjoint on each  $A_s(i)$  that acts as:

$$^*:\mathcal{A}_s(i)\to\mathcal{A}_s(i),$$

For  $T(u) \in \{L_n(u), G_n(u)\}$  establishing **diagonalisability** amounts to:

$$\left(\begin{pmatrix} 1 & -u & \cdots & u \\ 1 & -u & \cdots & u \\ u & -u & \cdots & u \end{pmatrix}^* = \begin{pmatrix} 1 & -u & \cdots & u \\ 1 & \cdots & u \\ u & -u & \cdots & u \end{pmatrix}, \quad \left(\begin{array}{c} -u & -u & \cdots & u \\ 1 & \cdots & u \\ -u & -u & \cdots & u \\ -u & -u & \cdots & u \end{array}\right)^* = \begin{array}{c} -u & -u & \cdots & u \\ -u & -u & \cdots & u \\ -u & -u & \cdots & u \\ -u & -u & \cdots & u \end{array}$$

which is typically a property of Yang-Baxter integrable R-operators.

#### Theorem (XP, Rasmussen '23)

For FC, BMW and Liu planar algebras if T(u) is **Yang-Baxter integrable** then it is **polynomially integrable**.

The integrable structure of both local and global operators for such planar algebras is generated by a single spectral-independent operator!

Xavier Poncini

A planar-algebraic universe

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$$\left(\mathbf{M}\right)^{*} = \mathbf{M}$$

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For FC, BMW and Liu planar algebras if T(u) is **Yang-Baxter integrable** then it is **polynomially integrable**.

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A planar-algebraic universe

# Outlook

#### Outlook

Summary:

- Discrete conformal nets
- Planar algebras provide (almost) examples
- Continuity of representations remains a challenge
- A class of local and global operators are polynomially integrable

Future work:

- Continuity of representations for all of  $T_k$
- More local and global operators
- Continuum limit taking discrete conformal nets to conformal nets
- Generalise to discrete algebraic quantum field theory

The end!